## Construction of the $\mathrm{C}_{2 \mathrm{v}}$ character table

The character table $\mathrm{C}_{2 \mathrm{v}}$ has the following form:

| $\mathrm{C}_{2 \mathrm{v}}$ | E | $\mathrm{C}_{2}$ | $\sigma_{\mathrm{v}}(\mathrm{xz})$ | $\sigma_{\mathrm{v}}{ }^{\prime}(\mathrm{yz})$ |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathrm{A}_{1}$ | 1 | 1 | 1 | 1 | z | $\mathrm{x}^{2}, \mathrm{y}^{2}, \mathrm{z}^{2}$ |
| $\mathrm{~A}_{2}$ | 1 | 1 | -1 | -1 | $\mathrm{R}_{\mathrm{z}}$ | xy |
| $\mathrm{B}_{1}$ | 1 | -1 | 1 | -1 | $\mathrm{x}, \mathrm{R}_{\mathrm{y}}$ | xz |
| $\mathrm{B}_{2}$ | 1 | -1 | -1 | 1 | $\mathrm{y}, \mathrm{R}_{\mathrm{x}}$ | yz |

We will consider how the various basis functions $x, y, z$ and others in the right hand columns map onto the basis Vectors of the space (also known as irreducible representations). These are the $\mathrm{A}_{1}, \mathrm{~A}_{2}, \mathrm{~B}_{1}$ and $\mathrm{B}_{2}$.

## The totally symmetric basis vector

As an example we will consider the $\mathrm{C}_{2 \mathrm{v}}$ point group, which corresponds to the $\mathrm{H}_{2} \mathrm{O}$ molecule for instance. We can call the totally symmetry basis vector $\mathrm{A}_{1}$ as shown in Table 1. $A_{1}$ is just a name like, $X, Y$ or $Z$ for the Cartesian space.

| $\mathbf{C}_{2 v}$ | $E$ | $\mathrm{C}_{2}$ | $\sigma_{\mathrm{v}(\mathrm{zz})}$ | $\sigma_{\mathrm{v}(y \mathrm{z})}$ |
| :--- | :--- | :--- | :--- | :--- |
| $\mathrm{A}_{1}$ | 1 | 1 | 1 | 1 |

The basis vectors must be normalized and orthogonal. Normalization is like saying that these are unit vectors in the space. The dimensionality of the basis must also equal the dimensionality of the space. The $\mathrm{C}_{2 \mathrm{v}}$ point group is 4-dimensional (i.e. there are four symmetry elements). To determine the dimensionality of any point group we need only count the symmetry operations.

## The normalization condition

Is $A_{1}$ normalized (i.e. does it have a length of 1 )? We will sum over the square of the contribution for each symmetry element and then divide by the dimension of the group $h$. The length $L$ is

$$
|S|^{2}=\frac{1}{\mathrm{~h}} \sum_{\mathrm{i}=1}^{\mathrm{h}} \chi_{\mathrm{i}}^{2}
$$

where $S$ is the basis vector name and $\chi$ is the character or value of each symmetry operation for that particular basis vector. We can see that $A_{1}$ is normalized.

$$
\left|A_{1}\right|^{2}=\frac{1}{4}\left(1^{2}+1^{2}+1^{2}+1^{2}\right)=1
$$

The other basis vectors must be orthogonal to $\mathrm{A}_{1}$ and also normalized. We could find these vectors using four equations and four unknowns.

## Starting point for construction

We will call the remaining vectors $A_{2}, B_{1}$ and $B_{2}$. Based on the information we have up to now we can construct a table of the basis vectors that looks like this.

| $\mathrm{C}_{2 \mathrm{v}}$ | E | $\mathrm{C}_{2}$ | $\sigma_{\mathrm{v}(\mathrm{yz)}}$ | $\sigma_{\mathrm{v}(\mathrm{yz})}$ |
| :--- | :--- | :--- | :--- | :--- |
| $\mathrm{A}_{1}$ | 1 | 1 | 1 | 1 |
| $\mathrm{~A}_{2}$ | 1 |  |  |  |
| $\mathrm{~B}_{1}$ | 1 |  |  |  |
| $\mathrm{~B}_{2}$ | 1 |  |  |  |

One formal way to find the remaining elements of the table is to use 9 unknown values that fill the table and then to set up nine equations based normalization, orthogonality and the sum rule that each column (except E) sums to zero. However, we will invoke an intuitive approach to find the basis vectors.

## The normalization condition for $A_{2}$

The subscript refers to whether a basis vector changes sign upon reflection. The letter describes whether it changes sign upon rotation. A does not change upon rotation, but B does. 1 does not change sign upon reflection, but 2 does.
Thus, $A_{2}$ does not change sign upon rotation ( $C_{2}=1$ ), but it does change sign upon reflection, i.e $\sigma_{v(x z)}=\sigma_{v(y z)}=-1$

| $\mathrm{C}_{2 \mathrm{v}}$ | E | $\mathrm{C}_{2}$ | $\sigma_{\mathrm{v}(x z)}$ | $\sigma_{\mathrm{v}(\mathrm{yz})}$ |
| :--- | :--- | :--- | :--- | :--- |
| $\mathrm{A}_{1}$ | 1 | 1 | 1 | 1 |
| $\mathrm{~A}_{2}$ | 1 | 1 | -1 | -1 |
| $\mathrm{~B}_{1}$ | 1 |  |  |  |
| $\mathrm{~B}_{2}$ | 1 |  |  |  |

We can see that $A_{2}$ is normalized sine

$$
\left|A_{2}\right|^{2}=\frac{1}{4}\left(1^{2}+1^{2}+(-1)^{2}+(-1)^{2}\right)=1
$$

The orthogonality is given by

$$
\left|\mathrm{S}_{1}\right|\left|\mathrm{S}_{2}\right|=\frac{1}{\mathrm{~h}} \sum_{\mathrm{i}=1}^{\mathrm{h}} \chi_{1 \mathrm{i}} \chi_{2 \mathrm{i}}
$$

which is explicitly given by

$$
\left|A_{1}\right|\left|A_{2}\right|=\frac{1}{4}((1)(1)+(1)(1)+(1)(-1)+(1)(-1))=0
$$

Next we consider the B basis vectors, which should change sign upon rotation.

| $\mathbf{C}_{2 \mathrm{v}}$ | $E$ | $\mathrm{C}_{2}$ | $\sigma_{\mathrm{v}(\mathrm{xz)}}$ | $\sigma_{\mathrm{v}(\mathrm{yz})}$ |
| :--- | :--- | :--- | :--- | :--- |
| $\mathrm{A}_{1}$ | 1 | 1 | 1 | 1 |
| $\mathrm{~A}_{2}$ | 1 | 1 | -1 | -1 |
| $\mathrm{~B}_{1}$ | 1 | -1 |  |  |
| $\mathrm{~B}_{2}$ | 1 | -1 |  |  |

## The orthonormal basis in $\mathrm{C}_{2 \mathrm{v}}$

| $\mathbf{C}_{\mathbf{2 v}}$ | $E$ | $C_{2}$ | $\sigma_{v(x z)}$ | $\sigma_{v(y z)}$ |
| :--- | :--- | :--- | :--- | :--- |
| $A_{1}$ | 1 | 1 | 1 | 1 |
| $A_{2}$ | 1 | 1 | -1 | -1 |
| $B_{1}$ | 1 | -1 | 1 | -1 |
| $B_{2}$ | 1 | -1 | -1 | 1 |


| $\mathrm{C}_{2 \mathrm{v}}$ | E | $\mathrm{C}_{2}$ | $\sigma_{\mathrm{v}}$ | $\sigma_{\mathrm{v}}{ }^{\prime}$ |
| :---: | :---: | :---: | :---: | :---: |
| $\Gamma_{1}$ | 1 | 1 | 1 | 1 |
| $\Gamma_{2}$ | 1 | 1 | -1 | -1 |
| $\Gamma_{3}$ | 1 | -1 | 1 | -1 |
| $\Gamma_{4}$ | 1 | -1 | -1 | 1 |

Mulliken symbols for irreps:

- "A"—symmetric wrt rotation about principle axis

$$
\left(\chi\left[C_{n}(z)\right]=+1\right)
$$

- "B"- irrep is antisymmetric wrt rotation about the principle axis $\left(\chi\left[\mathrm{C}_{n}(\mathrm{z})\right]=-1\right)$
- "E" -doubly degenerate irrep $(\mathrm{d}=2 \Rightarrow \chi(\mathrm{E})=2)$
- "T"-triply degenerate irrep $(\mathrm{d}=3 \Rightarrow \chi(\mathrm{E})=3)$


## Construction of the $\mathrm{C}_{3 \mathrm{v}}$ character table

The character tables $\mathrm{C}_{3 \mathrm{v}}$ has the following form:

| $\mathrm{C}_{3 \mathrm{v}}$ | E | $2 \mathrm{C}_{3}$ | $3 \sigma_{\mathrm{v}}$ |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathrm{A}_{1}$ | 1 | 1 | 1 | $z$ | $\mathrm{x}^{2}+\mathrm{y}^{2} ; \mathrm{z}^{2}$ |
| $\mathrm{~A}_{2}$ | 1 | 1 | -1 | $\mathrm{R}_{\mathrm{z}}$ |  |
| E | 2 | -1 | 0 | $(\mathrm{x}, \mathrm{y}) ;\left(\mathrm{R}_{\mathrm{x}}, \mathrm{R}_{\mathrm{y}}\right)$ | $\left(\mathrm{x}^{2}-\mathrm{y}^{2}, \mathrm{xy}\right) ;(\mathrm{xz}, \mathrm{yz})$ |
|  |  |  |  |  |  |

This character table contains a two-dimensional representation called $E$. This is because the $x$ and $y$ dimensions are coupled in this point group.

Let's generate the $\mathrm{C}_{3 \mathrm{v}}$ point group.
The operations of $\mathrm{C}_{3 \mathrm{v}}$ are $\mathrm{E}, 2 \mathrm{C}_{3}, 3 \sigma_{\mathrm{v}}(\mathrm{h}=6, \mathrm{~m}=3)$
$\mathrm{d}_{1}{ }^{2}+\mathrm{d}_{2}{ }^{2}+\mathrm{d}_{3}{ }^{2}=6$
$d_{1}=d_{2}=1$ and $d_{3}=2$.
Since the dimensions of the irreps are the $\chi(E)$ and every group contains the totally symmetric irrep,

| $\mathrm{C}_{3 \mathrm{v}}$ | $\mathbf{1 E}$ | $\mathbf{2} \mathrm{C}_{3}$ | $\mathbf{3} \sigma_{\mathrm{v}}$ |
| :---: | :---: | :---: | :---: |
| $\Gamma_{1}$ | 1 | 1 | 1 |
| $\Gamma_{2}$ | 1 | j | k |
| $\Gamma_{3}$ | 2 | m | n |

Orthogonality of $\Gamma_{1} \& \Gamma_{2}:(1)(1)(1)+(2)(1)(j)+(3)(1)(k)=0$ $1+2 j+3 k=0$
Normalization of $\Gamma_{2}$ requires (1)(1) ${ }^{2}+2(\mathrm{j})^{2}+3(\mathrm{k})^{2}=6$ $\mathrm{j}=+1$ and $\mathrm{k}=-1$
Orthogonality of $\Gamma_{1} \& \Gamma_{3}:(1)(1)(2)+(2)(1)(m)+(3)(1)(n)=0$
$2+2 m+3 n=0$
Normalization of $\Gamma_{3}$ means $(1)(2)^{2}+2(m)^{2}+3(n)^{2}=6$ $\mathrm{m}=-1$ and $\mathrm{n}=0$.

| $\mathrm{C}_{3 \mathrm{v}}$ | $\mathbf{1 E}$ | $\mathbf{2 \mathrm { C } _ { 3 }}$ | $\mathbf{3} \sigma_{\mathrm{v}}$ |
| :---: | :---: | :---: | :---: |
| $\Gamma_{1}$ | 1 | 1 | 1 |
| $\Gamma_{2}$ | 1 | j | k |
| $\Gamma_{3}$ | 2 | m | n |

## Mulliken symbols for $\mathrm{C}_{3 \mathrm{v}}$

| $\mathrm{C}_{3 \mathrm{v}}$ | $\mathbf{1 E}$ | $\mathbf{2 C _ { 3 }}$ | $\mathbf{3} \sigma_{\mathrm{v}}$ |
| :---: | :---: | :---: | :---: |
| $\mathrm{A}_{1}$ | 1 | 1 | 1 |
| $\mathrm{~A}_{2}$ | 1 | 1 | -1 |
| E | 2 | -1 | 0 |

Treating rotations and binary products as before, we can represent the C3v point group as

| $\mathrm{C}_{3 \mathrm{v}}$ | E | $2 \mathrm{C}_{3}$ | $3 \sigma_{\mathrm{v}}$ |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathrm{A}_{1}$ | 1 | 1 | 1 | $z$ | $\mathrm{x}^{2}+\mathrm{y}^{2} ; \mathrm{z}^{2}$ |
| $\mathrm{~A}_{2}$ | 1 | 1 | -1 | $\mathrm{R}_{\mathrm{z}}$ |  |
| E | 2 | -1 | 0 | $(\mathrm{x}, \mathrm{y}) ;\left(\mathrm{R}_{\mathrm{x}}, \mathrm{R}_{\mathrm{y}}\right)$ | $\left(\mathrm{x}^{2}-\mathrm{y}^{2}, \mathrm{xy}\right) ;(\mathrm{xz}, \mathrm{yz})$ |

The $x 2-y 2$ and $x y$ orbitals are also degenerate as are the $x z$ and $y z$ orbitals

In many instances there are more than one $A, B, E$, etc. irreps present in the point group so subscripts and superscripts are used
$\boldsymbol{g}$ or $\boldsymbol{u}$ subscripts are used in point groups with centers of symmetry (i) to denote gerade (symmetric) and ungerade (antisymmetric) with respect to inversion
' and " are used to designate symmetric and antisymmetric with respect to inversion through a $\sigma_{h}$ plane
numerical subscripts are used otherwise

